

B.

M.H.DBennet's relation:-

We consider an infinite cylinder of conductive fluid with an axial current density  $J_z = J_z(r)$  and a resultant of azimuthal magnetic induction  $B_\theta = B_\theta(r)$ . For simplicity the current density, the magnetic field pressure etc are assumed to depend on the distance from the cylinder's axis.

Here viscous and gravitational forces are neglected. Thus we have eqn of motion as -

$$0 = -\vec{\nabla} p + \frac{1}{c} (\vec{J} \times \vec{B}) \quad \text{--- (1)}$$

using cylindrical co-ordinates the eqn (1) can be expressed as -

$$0 = -\frac{dp}{dr} - \frac{d}{dr} \left( \frac{B_\theta^2}{8\pi} \right) - \frac{B_\theta^2}{4\pi r}$$

$$\Rightarrow \frac{dp}{dr} = -\frac{1}{8\pi r^2} \frac{d}{dr} (r^2 B_\theta^2)$$

From Maxwell's eqns ~~we~~ we have the following relation

$$\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{J} \quad \text{--- (2)}$$

using cylindrical co-ordinates eqn (2) can be expressed as -

$$\frac{1}{r} \frac{d}{dr} (r B_\theta) = \frac{4\pi}{c} J_z(r)$$

$$\Rightarrow \frac{d}{dr} (r B_\theta) = \frac{4\pi r}{c} J_z(r) \quad \text{--- (3)}$$

integrating (1) we get -

$$p_0 = \frac{4\pi}{c_s} \int_0^R r J_2(r) dr$$

$$= \frac{2I}{c_s} \quad \text{--- (5)}$$

where  $I = 2\pi \int_0^R r J_2(r) dr$

Integrating (2) we name the solution as

$$p = p_0 - \frac{1}{8\pi} \int_0^R \frac{1}{r^2} \frac{d}{dr} (r^2 \rho a^2) dr \quad \text{--- (6)}$$

where  $p_0$  is the pressure of the fluid at  $r=0$  that is on  $z$  axis. If the pressure is continued to  $r \leq R$ , the pressure drops to zero at  $r=R$ . Consequently, the axial pressure  $p_0$  is given by

$$p_0 = \frac{1}{8\pi} \int_0^R \frac{1}{r^2} \frac{d}{dr} (r^2 \rho a^2) dr \quad \left. \begin{array}{l} r=0 \\ p=0 \end{array} \right\}$$

using (7) from eqn (6) we get -

$$p = \frac{1}{8\pi} \int_r^R \frac{1}{r^2} \frac{d}{dr} (r^2 \rho a^2) dr \quad \text{--- (8)}$$

The average pressure  $\langle P \rangle$  can be defined as -

$$\langle P \rangle = \frac{1}{\pi R^2} \int_0^R (2\pi r) p dr \quad \text{--- (9)}$$

integrating by parts we have -

$$\langle P \rangle = \frac{1}{\pi R^2} \left[ 2\pi \left\{ \left( \frac{\rho a^2}{2} \right)_0^R - \int_0^R \frac{dp}{dr} \cdot \frac{r^2}{2} dr \right\} \right]$$

$$\begin{aligned}
&= \frac{2}{R^2} \left[ 0 - \frac{1}{2} \int_0^R r^2 \frac{dp}{dr} dr \right] \\
&= - \frac{1}{R^2} \int_0^R r^2 \frac{dp}{dr} dr \\
&= - \frac{1}{R^2} \int_0^R r^2 \left( - \frac{1}{8\pi r^2} \frac{d}{dr} (r^2 B^2) \right) dr \\
&\quad \text{(using ②)} \\
&= \frac{1}{8\pi R^2} \int_0^R \frac{d}{dr} \left( r^2 \cdot \frac{4\mu^2}{c^2 r^2} \right) dr \\
&= \frac{1}{8\pi R^2} \cdot \frac{4\mu^2}{c^2} \int_0^R \frac{d}{dr} \left( \frac{4\mu^2}{c^2} \right) dr \\
&\quad \text{using ⑤} \\
&= \frac{1}{8\pi R^2} \cdot \frac{4\mu^2}{c^2} \\
&= \frac{\mu^2}{2\pi R^2 c^2} \quad \text{⑩}
\end{aligned}$$

Note that the average pressure  $\langle p \rangle$  of the matter is equal to the magnetic pressure  $\frac{B^2}{8\pi}$  at the surface of the cylinder.

If the electrons and ions are at the same temperature  $T$  and number density of electron per unit length  $N_e$ , then

$N_e$  can be expressed as

$$N_e = \int_0^R (2\pi r) N dr$$

where  $N$  is the total number density and can be expressed as

$$N = \frac{P}{2kT}$$

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$$N_e = \int_0^R (2\pi r) \cdot \frac{P}{2kT} dr$$

$$= \frac{2\pi}{kT} \int_0^R r P dr$$

$$N_e = \frac{\pi}{kT} \left\{ \left( \frac{r^2}{2} P \right)_0^R - \frac{1}{2} \int_0^R \frac{dP}{dr} r^2 dr \right\}$$

$$= \frac{\pi}{2kT} \int_0^R r^2 \cdot \left( -\frac{1}{8\pi r^2} \frac{d}{dr} (r^2 B_0^2) dr \right)$$

$$= + \frac{\pi \cdot 1}{16kT} \int_0^R \frac{d}{dr} (r^2 B_0^2) dr$$

$$= \frac{1}{16kT} \int_0^R \frac{d}{dr} \left( r^2 \frac{4I^2}{c^2 r^2} \right) dr$$

$$= \frac{1}{16kT} \frac{4I^2}{c^2}$$

$$N_e = \frac{I^2}{4kTc^2}$$

$$I^2 = 4kTc^2 N_e \quad \text{--- (11)}$$

This is the Bennett's relation for linear pinch. It shows that the kinetic temperature  $T$  is proportional to the square of the total current  $I$  flowing in the discharge tube and

inversely to the wave density  $N_e$  to the particle.

$$T = \left( \frac{1}{2k^2} \right) \frac{I^2}{N_e}$$

Def.<sup>n</sup> (Note)

Derivation of Generalised OHM'S LAW

Let us define the mean velocity and mass <sup>density</sup> velocity  $\rho_m$  of current density  $J$  and charge density  $\rho_q$

$$v = \frac{m_i n_i v_i + m_e n_e v_e}{m_i n_i + m_e n_e} \quad (1)$$

$$J = e (m_i n_i v_i - m_e n_e v_e) \quad (2)$$

$$\rho_m = m_i n_i + m_e n_e \quad (3)$$

$$\rho_q = e (n_i - n_e) \quad (4)$$

Now the momentum transfer eqns for electrons and ions are

$$(1) \quad \frac{\partial \vec{v}_e}{\partial t} + (\vec{v}_e \cdot \nabla) \vec{v}_e = - \frac{\nabla P_e}{m_e n_e} - \frac{e}{m_e} \left[ \vec{E} + \frac{1}{c} (\vec{v}_e \times \vec{B}_0) \right]$$

and collision term:

$$\frac{\partial \vec{v}_i}{\partial t} + (\vec{v}_i \cdot \nabla) \vec{v}_i = - \frac{\nabla P_i}{m_i n_i} + \frac{e}{m_i} \left[ \vec{E} + \frac{1}{c} (\vec{v}_i \times \vec{B}_0) \right] + \text{collision term}$$

total current of the total current

Multiplying equation (5) by  $-e n_e$  and (6) by  $e n_i$  and then adding we get -

$$\begin{aligned}
 & -e n_e \frac{\partial v_e}{\partial t} + e n_i \frac{\partial v_i}{\partial t} - e n_e (v_e \cdot \nabla) v_e \\
 & + e n_i (v_i \cdot \nabla) v_i = e \left( \frac{\nabla p_e}{m_e} - \frac{\nabla p_i}{m_i} \right) \\
 & + \left( \frac{e^2 n_e}{m_e} + \frac{e^2 n_i}{m_i} \right) E + \frac{1}{c} \left( \frac{e^2 n_e v_e}{m_e} + \frac{e^2 n_i v_i}{m_i} \right) \times B_0
 \end{aligned}$$

We have the equation of continuity for electrons & ions as

$$\frac{\partial n_e}{\partial t} + \nabla \cdot (n_e v_e) = 0 \quad \text{--- (8)}$$

$$\frac{\partial n_i}{\partial t} + \nabla \cdot (n_i v_i) = 0 \quad \text{--- (9)}$$

Now

$$e n_i \frac{\partial v_i}{\partial t} - e n_e \frac{\partial v_e}{\partial t} = \frac{\partial}{\partial t} (e n_i v_i - e n_e v_e)$$

$$\begin{aligned}
 & = \frac{\partial j}{\partial t} + e v_i \nabla (n_i v_i) + v_e \frac{\partial}{\partial t} (e n_e) \\
 & \quad - e v_e \nabla (n_e v_e) \quad \text{(using 2)}
 \end{aligned}$$

now

$$\frac{e^2 n_e v_e}{m_e} = \frac{e}{m_e} (e n_e v_e - e n_i v_i)$$

$$+ \frac{e^2 n_i v_i}{m_e}$$

$$= \frac{e^2 n_i v_i}{m_e} - \frac{e}{m_e} j \quad \text{--- (10)}$$